The Fokker-Planck Equation

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$$d\mathbf{y}_t = \mathbf{u}(\mathbf{y}_t, t) dt + \mathbf{L}(\mathbf{y}_t, t) d\mathbf{W}_t, \quad \mathbf{y}_0 \sim \rho(\mathbf{y}_0, 0)$$

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where $\boldsymbol{u} \colon \mathbb{R}^n \times [0,\infty) \to \mathbb{R}^n$ is the drift, $\boldsymbol{L} \colon \mathbb{R}^n \times [0,\infty) \to \mathbb{R}^{n \times m}$ is the diffusion, and \boldsymbol{W}_t is *m*-dimensional Brownian noise.

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where

$$\int_0^{\tau} \boldsymbol{L}(\boldsymbol{y}_{\tau},\tau) \,\mathrm{d}\boldsymbol{W}_t = \lim_{\delta t \to 0} \sum_{[\tau_i,\tau_{i+1}] \in \mathcal{P}_{[0,t]}} \boldsymbol{L}(X_{\tau_i},\tau_i) \left(\boldsymbol{W}\tau_{i+1} - \boldsymbol{W}_{\tau_i}\right).$$

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This can be thought of as a generalisation of a Riemann-Stieltjes integral.

$$d\mathbf{y}_t = \mathbf{u}(\mathbf{y}_t, t) dt + \mathbf{L}(\mathbf{y}_t, t) d\mathbf{W}_t.$$

• The solution y_t is a stochastic process taking values in \mathbb{R}^n .

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- In general, an SDE cannot be solved analytically...
- ...but we can solve it numerically to obtain sample paths.

e.g.
$$dX_t = X_t dt + dW_t$$











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e.g. $dX_t = tanh(X_t) dt + dW_t$

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The probability density function $\rho(\mathbf{x}, t)$ of \mathbf{y}_t at time t satisfies the Fokker-Planck equation

$$\frac{\partial \rho(\boldsymbol{x},t)}{\partial t} = -\boldsymbol{\nabla} \cdot \left(\rho(\boldsymbol{x},t) \, \boldsymbol{u}(\boldsymbol{x},t) \right) + \frac{1}{2} \boldsymbol{\nabla} \cdot \boldsymbol{\nabla} \cdot \left(\rho(\boldsymbol{x},t) \, \boldsymbol{L}(\boldsymbol{x},t) \, \boldsymbol{L}(\boldsymbol{x},t)^{\mathsf{T}} \right)$$

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with the additional constraints

$$\int_{\mathbb{R}^n}
ho({m{x}},t) \, \mathsf{d}{m{x}} = 1, \quad \lim_{\|{m{x}}\| o \infty}
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This presents an alternative to solving the SDE directly.

$$\frac{\partial \rho(\mathbf{x},t)}{\partial t} = -\nabla \cdot \left(\rho(\mathbf{x},t) \, \mathbf{u}(\mathbf{x},t) \right) + \frac{1}{2} \nabla \cdot \nabla \cdot \left(\rho(\mathbf{x},t) \, \mathbf{L}(\mathbf{x},t) \, \mathbf{L}(\mathbf{x},t)^{\mathsf{T}} \right)$$

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Provide better insight into rare events and distribution tails.

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- Provide an easier way to incorporate boundary conditions.
- Provide insight into the long-term behaviour of a system we can find a steady state solution.

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- Provide better insight into rare events and distribution tails.
- Provide an easier way to incorporate boundary conditions.
- Provide insight into the long-term behaviour of a system we can find a steady state solution.
- Provide a way to map uncertain initial conditions forward with SDE dynamics — this is useful in data assimilation.

$$dy_t = dW_t, \quad y_0 = 0 \tag{1}$$

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$$\frac{\partial \rho(x,t)}{\partial t} = \frac{1}{2} \nabla^2 \rho(x,t), \quad \rho(x,0) \ "=" \ \delta(x)$$

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This is exactly the heat equation. The solution is $y_t \sim \mathcal{N}(0, t)$ and so

$$\rho(x,t) = \frac{1}{t\sqrt{2\pi}} \exp\left[-\frac{x^2}{2t}\right]$$



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Now, consider

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The solution is a diffusing Gaussian with a drifting mean

$$\rho(x,t) = \frac{1}{t\sqrt{2\pi}} \exp\left[-\frac{\left(x-e^{t}\right)^{2}}{2t}\right]$$

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 c is a scalar field of interest (e.g. temperature, concentration of a pollutant, density of an organism),

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The advection-diffusion PDE is

$$\frac{\partial c(\boldsymbol{x},t)}{\partial t} = \nabla \cdot \left(D \nabla c(\boldsymbol{x},t) - \boldsymbol{v}(\boldsymbol{x},t) c(\boldsymbol{x},t) \right)$$

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This means that every advection-diffusion problem can be equivalently written as a stochastic differential equation:

$$\mathrm{d}\boldsymbol{x}_t = \boldsymbol{v}(\boldsymbol{x}_t, t) \, \mathrm{d}t + D \, \mathrm{d}\boldsymbol{W}_t.$$

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Consider quantifying transport of material in the ocean.

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This data is only available at certain gridpoints, so it is missing features of the true flow, e.g. subgrid-scale eddies. One common approach is to model these as stochastic perturbations.

A simple example would be:

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$$d\begin{pmatrix} x_t \\ y_t \end{pmatrix} = \left[\begin{pmatrix} u(x_t, y_t, t) \\ v(x_t, y_t, t) \end{pmatrix} + \begin{pmatrix} \mu_t \\ \nu_t \end{pmatrix} \right] dt$$
$$d\begin{pmatrix} \mu_t \\ \nu_t \end{pmatrix} = -\frac{1}{T_L} \begin{pmatrix} \mu_t \\ \nu_t \end{pmatrix} dt + \kappa(x_t, y_t) d\begin{pmatrix} W_t^{(1)} \\ W_t^{(2)} \end{pmatrix}$$

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In our notation

• The state variable is $\mathbf{x} = (x_t, y_t, \mu_t, \nu_t)$, so n = 4 and m = 2. The velocity is

$$\boldsymbol{u}(\boldsymbol{x},t) = \begin{pmatrix} u(x_t,y_t,t) \\ v(x_t,y_t,t) \\ -\frac{1}{T_L}\mu_t \\ -\frac{1}{T_L}\nu_t \end{pmatrix}$$

 \blacktriangleright The diffusion is the 4 \times 2 matrix

$$\boldsymbol{L}(\boldsymbol{x},t) = \begin{pmatrix} \boldsymbol{O}_{2\times 2} \\ \boldsymbol{\kappa}(x_t,y_t) \end{pmatrix}$$





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t = 4 days





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